# The SUSY-Yukawa sum rule

#### Testing weak scale naturalness at the LHC

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Brookhaven Forum 2010: A space-time odyssey BNL – May 26, 2010

#### Goals for the next 15 minutes

- introduce SUSY-Yukawa sum rule that allows to test supersymmetric stabilisation of the weak scale
- Denchmark study for LHC prospects to test the sum rule
  - here: focus on lighter stop and sbottom masses
    - > can be precisely determined at the LHC
  - progress on stop mixing possible
  - sbottom mixing and heavier masses difficult

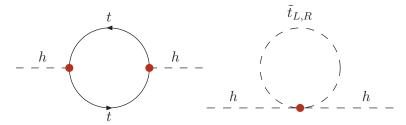




MB, D. Curtin, M. Perelstein, 1004.5350

# Introduction: SUSY cancellation of quadratic divergences

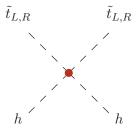
• hierarchy problem: loop contributions of SM particles (e.g. tops) let the Higgs potential depend quadratically on the cut-off scale



- new particles (stops) with sub-TeV masses required to cancel these contributions
- couplings to the Higgs boson have to be equal

# How to access the stop-Higgs coupling at the LHC?

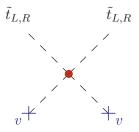
We want to measure the coupling  $hh ilde{t}_{L,R} ilde{t}^c_{L,R}$ :



direct measurement not feasible at the LHC

# How to access the stop-Higgs coupling at the LHC?

We want to measure the coupling  $hh ilde{t}_{L,R} ilde{t}^c_{L,R}$ :



- direct measurement not feasible at the LHC
- EWSB ➤ contribution to stop mass matrix

# The stop mass matrix

• stop mass matrix:  $\mathcal{L} = (\tilde{t}_L^c, \tilde{t}_R^c) \mathcal{M}_t^2(\tilde{t}_L, \tilde{t}_R)$ 

$$\mathcal{M}_t^2 = \begin{pmatrix} m_L^2 + m_t^2 + \Delta_u & m_t(A_t + \mu \cot \beta) \\ m_t(A_t + \mu \cot \beta) & m_R^2 + m_t^2 + \Delta_{\bar{u}} \end{pmatrix}$$

rotation to mass eigenstates via

$$\begin{array}{lll} \tilde{t}_1 & = & \cos \theta_t \, \tilde{t}_L + \sin \theta_t \, \tilde{t}_R \\ \tilde{t}_2 & = & -\sin \theta_t \, \tilde{t}_L + \cos \theta_t \, \tilde{t}_R \end{array}$$

ullet then re-express  $\mathcal{M}_{11}^2$ 

$$m_L^2 + \frac{m_t^2}{m_t^2} + \Delta_u = m_{\tilde{t}_1}^2 \cos^2 \theta_t + m_{\tilde{t}_2}^2 \sin^2 \theta_t$$

analogously for sbottom system

#### The SUSY-Yukawa sum rule

eliminating  $m_L^2$  yields the SUSY-Yukawa sum rule  $(m_b \rightarrow 0)$ 

$$m_t^2 + (\Delta_u - \Delta_d) = m_{\tilde{t}_1}^2 \cos^2 \theta_t + m_{\tilde{t}_2}^2 \sin^2 \theta_t - m_{\tilde{b}_1}^2 \cos^2 \theta_b - m_{\tilde{b}_2}^2 \sin^2 \theta_b$$

where 
$$\Delta_u - \Delta_d = m_Z^2 \cos^2 \theta_W \cos 2\beta \approx -m_W^2$$

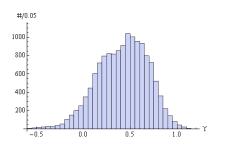
- sum rule expresses stop-Higgs coupling in terms of measurable quantities (masses, mixing angles)
- SUSY weak scale stabilization (in principle) testable at the LHC!

#### ... and what about radiative corrections?

- above derivation valid at tree level
- to quantify effect of radiative corrections, define

$$\Upsilon = \frac{1}{v^2} (m_{\tilde{t}_1}^2 \cos^2 \theta_t + m_{\tilde{t}_2}^2 \sin^2 \theta_t - m_{\tilde{b}_1}^2 \cos^2 \theta_b - m_{\tilde{b}_2}^2 \sin^2 \theta_b)$$

• SUSY tree level prediction:  $\Upsilon_{\text{tree}} = 0.28 \pmod{\beta}$  a few)



 SuSpect scan over pMSSM parameter space yields

$$0 \lesssim \Upsilon \lesssim 1$$

'generic' prediction:  $|\Upsilon| < 16\pi^2$ 

 can be narrowed by measuring some SUSY masses (see later)

#### Parameters to be determined

masses

 $m_{ ilde{t}_1}, m_{ ilde{t}_2} \ m_{ ilde{b}_1}, m_{ ilde{b}_2}$ 

mixing angles

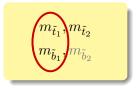
 $\sin heta_t \\ \sin heta_b \ ext{(usually small)}$ 

 $\tan \beta$  (minor impact)

additional information helpful (to pin down radiative corrections)

#### Parameters to be determined

masses



mixing angles

 $\sin \theta_t$   $\sin \theta_b$  (usually small)  $\tan \beta$  (minor impact)

additional information helpful (to pin down radiative corrections)

# Our benchmark scenario - main virtues

ightharpoonup study  $pp o ilde{g} ilde{g} o bb ilde{b} ilde{b}_1 ilde{b}_1 o 4b+
ot\!\!\!/ _T$  and  $pp o ilde{t}_1 ilde{t}_1^c o tar{t}+
ot\!\!\!/ _T$ 

$$\tilde{g} \stackrel{b}{\underset{b}{\tilde{b}_{1}}} \tilde{b}_{1} \tilde{\chi}_{1}^{0}$$



(parton level analysis using MG/ME and BRIDGE)

# Gluino pair production – $4b + E_T$

#### It's all about finding edges!

- ullet consider  $pp o 2 ilde{g} o 2b + 2 ilde{b}_1 o 4b + E_T$
- $\sigma(pp o 2\tilde{g}) \simeq 11 \, \mathrm{pb} \, \, \mathrm{for} \, \sqrt{s} = 14 \, \mathrm{TeV}$
- with  $\mathcal{L} = 10 \, \text{fb}^{-1}$ :  $\sim 4800 \, \text{signal events}$ , SM background negligible!



however: combinatorial background – which b-jet is which?

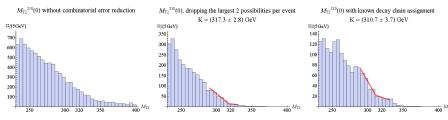
several possible ways to get rid of 'wrong' pairings (e.g.  $\Delta R$ )

> we always require **two independent methods** to yield **consistent results** (otherwise measurement is rejected)

# Mass determination for $\tilde{g}$ , $\tilde{b}_1$ and $\tilde{\chi}^0_1$

- ullet  $M_{bb}$  invariant mass endpoint can easily be recovered
- need two more edges to pin down all masses  $\succ M_{T2}^{(2,2,0)}$ ,  $M_{T2}^{(2,1,0)}$

Barr et al, hep-ph/0304226; Burns et al, 0810.5576



> combining those, we obtain the mass measurements

	68% C.L.	theory
$\overline{m_{\tilde{b}_1}}$	(316,356)	341 GeV
$m_{ ilde{g}}$	(508,552)	525 GeV
$m_{ ilde{\chi}^0_1}$	$(45^{(*)},115)$	98 GeV

(\*) LEP lower bound

# Stop pair production – extracting $m_{ ilde{t}_1}$

- ullet analyze  $pp o 2 ilde{t}_1 o 2t+
  ot\!\!\!/ t_T$
- $\sigma(pp o 2 \tilde{t}_1) \simeq 2 \ {
  m pb} \ {
  m for} \ \sqrt{s} = 14 \ {
  m TeV}$
- impose standard cuts & use hadronic tops (following MEADE, REECE, HEP-PH/0601124)

 $\tilde{t}_1 = \tilde{\chi}_1^0 \\
\tilde{t}_1 = \tilde{\chi}_1^0$ 

- with  $\mathcal{L} = 100 \, \mathrm{fb}^{-1}$ :  $S/B \simeq 14$ ,  $S/\sqrt{B} = 140$
- straightforward to extract

$$(M_{T2})_{\rm max}(\chi=0) = (340 \pm 4) \, {\rm GeV}$$
 theory: 336.7 GeV

• using our previous  $m_{\tilde{\chi}^0_1}$  measurement we find (68% C.L.)

$$356\,\mathrm{GeV} \leq m_{\tilde{t}_1} \leq 414\,\mathrm{GeV}$$
 theory:  $371\,\mathrm{GeV}$ 

#### So what did we learn?

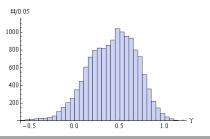
ullet rewrite  $\Upsilon$  as

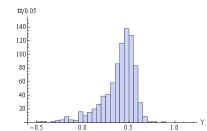
$$\Upsilon = \underbrace{\frac{1}{v^2}(m_{\tilde{t}_1}^2 - m_{\tilde{b}_1}^2)}_{\Upsilon'} + \underbrace{\frac{\sin^2 \theta_t}{v^2}(m_{\tilde{t}_2}^2 - m_{\tilde{t}_1}^2)}_{\Delta \Upsilon_t} - \underbrace{\frac{\sin^2 \theta_b}{v^2}(m_{\tilde{b}_2}^2 - m_{\tilde{b}_1}^2)}_{\Delta \Upsilon_b}$$

- our measurements yield
- $\Upsilon' = 0.53^{+0.20}_{-0.15}$

theory: 0.35

- ullet no information on  $\Delta \Upsilon_t, \Delta \Upsilon_b >$  lepton collider
- ullet however: much more accurate prediction for  $\Upsilon$



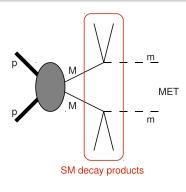


# Summary

- SUSY-Yukawa sum rule relates stop-Higgs coupling to stop and sbottom masses and mixing angles ➤ measurable quantities
- verifying (or falsifying) the sum rule means testing SUSY as the origin of the weak scale stabilization
- full measurement will have to wait for lepton collider
- We can make significant progress at the LHC in some regions of parameter space:
  - masses of  $\tilde{t}_1$ ,  $\tilde{b}_1$ ,  $\tilde{g}$  and  $\tilde{\chi}^0_1$  can be precisely measured  $\succ$  more accurate prediction for the sum rule
- $oldsymbol{\circ}$  we also developed new techniques to reduce combinatorial backgrounds in  $M_{T2}$  analyses

# Back-up slides

# The stransverse mass $M_{T2}$



Barr, Lester, Stephens, hep-ph/0304226

#### SUSY events complicated by

- two missing particles
- LSP mass unknown
- > the best we can do
  - use trial LSP mass  $\chi$
  - minimize over all possible LSP momentum configurations

 $\triangleright$  define the stransverse mass  $M_{T2}$  by

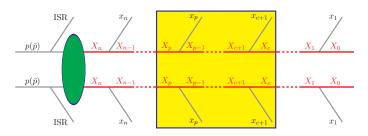
$$M_{T2}(\chi) = \min_{\boldsymbol{p}_T^{(1)} + \boldsymbol{p}_T^{(2)} = \boldsymbol{p}_T} \left\{ \max\{m_T^{(1)}, m_T^{(2)}\} \right\}$$

edge of distribution:  $M_{T2}(\chi)_{\rm max}=rac{M^2-m^2}{2M}+\sqrt{(rac{M^2-m^2}{2M})^2+\chi^2}$ 

# Extension: the subsystem $M_{T2}$

Burns, Kong, Matchev, Park, 0810.5576

#### for n > 1 step decay chains:



generalize  $M_{T2}$  concept to subsystem  $M_{T2}^{(n,p,c)}(\chi)$  (n: grandparent index, p: parent index, c: child index)

ightharpoons  $M_{T2}^{(n,p,c)}(\chi)$  endpoint yields relation between  $m_n$ ,  $m_p$  and  $m_c$ 

# Definition of benchmark scenario

parameter	EWSB scale value	
$M_1$	100 GeV	
$M_{2,3}$	450 GeV	
$A_t$	390 GeV	
$\mu$	400 GeV	
aneta	10	
$M_A$	600 GeV	
$m_{ ilde{e}_{L,R}, ilde{ au}_{L,R}, ilde{q}_L ilde{u}_R, ilde{d}_R}$	1000 GeV	
$m_{ ilde{Q}_L}$	310 GeV	
$m_{ ilde{t}_R}$	780 GeV	

# SM backgrounds to $4b + p_T$

Background	Generator	$\epsilon_{ m b}\sigma$	$\epsilon_{ m b}\epsilon_{ m kin}\sigma$
$4j + (Z \to \nu \nu)$	MGME, ALPGEN	10 fb	
diboson + jets	_	< 10 fb	
$tt \to n\tau + X$	MGME, BRIDGE	$21.6\mathrm{pb}$	25 fb
t	_		≪ 30 fb

assumed b-tagging efficiencies: 0.6 (b), 0.1  $(c, \tau)$ , 0.01 (light jet)

#### Some technical details

- SUSY spectrum and decays calculated using SUSY-HIT
- $\bullet$  parton-level analysis for  $\sqrt{s}=14\,\mathrm{TeV}~pp$  collisions
- Monte Carlo event samples generated by MadGraph/MadEvent
- fully decayed final state obtained with BRIDGE
- leading order analysis, using CTEQ6l1 pdf sets
- Gaussian smearing of jet energies

$$\frac{\Delta E}{E} = \frac{50\%}{\sqrt{E[\text{GeV}]}} \oplus 3\%$$

to simulate detector response